

Two-body problem

$$\vec{r} = \vec{r}_O + \vec{r}_{\text{rel}} \quad \vec{v} = \vec{v}_O + \vec{\Omega} \times \vec{r}_{\text{rel}} + \vec{v}_{\text{rel}} \quad \vec{a} = \vec{a}_O + \dot{\vec{\Omega}} \times \vec{r}_{\text{rel}} + \vec{\Omega} \times (\vec{\Omega} \times \vec{r}_{\text{rel}}) + 2\vec{\Omega} \times \vec{v}_{\text{rel}} + \vec{a}_{\text{rel}}$$

$$\ddot{\vec{r}} = -\frac{\mu}{r^3} \vec{r} \quad \vec{r}' = \vec{r}_2 - \vec{r}_1 \quad \vec{h} = \vec{r}' \times \vec{r} \quad h = v_{\perp} r = r^2 \dot{\theta}$$

$$r = \frac{h^2}{\mu} \frac{1}{1 + e \cos \theta} \quad \frac{dA}{dt} = \frac{h}{2} \quad \frac{v^2}{2} - \frac{\mu}{r} = \varepsilon = -\frac{\mu^2}{2h^2}(1 - e^2) \quad \mu = G(m_1 + m_2)$$

$$v_r = \frac{\mu}{h} e \sin \theta \quad v_{\perp} = \frac{h}{r} = \frac{\mu}{h}(1 + e \cos \theta) \quad \tan \gamma = \frac{v_r}{v_{\perp}} \quad r = R + z$$

Elliptical orbits ($0 \leq e < 1$)

$$2a = r_a + r_p \quad 2ae = r_a - r_p \quad h^2 = \mu a (1 - e^2) \quad \varepsilon = -\frac{\mu}{2a} \quad T^2 = \frac{4\pi^2}{\mu} a^3$$

$$b = a\sqrt{1 - e^2} = \sqrt{r_p r_a} \quad \bar{r}_{\theta} = a\sqrt{1 - e^2} = b \quad \bar{r}_t = a \left(1 + \frac{e^2}{2}\right)$$

$$\tan \frac{E}{2} = \sqrt{\frac{1 - e}{1 + e}} \tan \frac{\theta}{2} \quad M_e = E - e \sin E \quad M_e = \frac{\mu^2}{h^3} (1 - e^2)^{3/2} t = \frac{2\pi}{T} t$$

Parabolic orbits ($e = 1$)

$$v = \sqrt{\frac{2\mu}{r}} \quad \gamma = \frac{\theta}{2} \quad h^2 = 2\mu r_p$$

$$M_p = \frac{1}{2} \tan \frac{\theta}{2} + \frac{1}{6} \tan^3 \frac{\theta}{2} \quad M_e = \frac{\mu^2}{h^3} t \quad \tan \frac{\theta}{2} = \left[3M_p + \sqrt{(3M_p)^2 + 1}\right]^{1/3} - \left[3M_p + \sqrt{(3M_p)^2 + 1}\right]^{-1/3}$$

Hyperbolic orbits ($e > 1$)

$$\varepsilon = \frac{\mu}{2a} \quad v^2 - v_{\text{esc}}^2 = v_{\infty}^2 \quad v_{\infty} = \frac{\mu}{h} \sqrt{e^2 - 1} \quad 1 + e \cos \theta_{\infty} = 0 \quad h^2 = \mu a (e^2 - 1)$$

$$\tanh \frac{F}{2} = \sqrt{\frac{e - 1}{e + 1}} \tan \frac{\theta}{2} \quad M_h = e \sinh F - F \quad M_h = \frac{\mu^2}{h^3} (e^2 - 1)^{3/2} t$$

Newton-Raphson method

$$x_{i+1} = x_i - \frac{f(x_i)}{f'(x_i)}$$

Classical orbital elements

$$\begin{aligned}
 r &= \sqrt{\vec{r} \cdot \vec{r}} & v &= \sqrt{\vec{v} \cdot \vec{v}} & v_r &= \vec{v} \cdot \hat{r} & \vec{h} &= \vec{r} \times \vec{v} & h &= \sqrt{\vec{h} \cdot \vec{h}} & \cos i &= h_z/h \\
 \vec{N} &= \hat{K} \times \vec{h} \quad (N_x = -h_y, \quad N_y = h_x) & \cos \Omega &= N_x/N & \text{(if } N_y < 0 \text{ then } \Omega = 2\pi - \cos^{-1}) \\
 \vec{e} &= \frac{1}{\mu} \left[\left(v^2 - \frac{\mu}{r} \right) \vec{r} - r v_r \vec{v} \right] & e &= \sqrt{\vec{e} \cdot \vec{e}} & \cos \omega &= \vec{N} \cdot \vec{e}/Ne & \text{(if } e_z < 0 \text{ then } \omega = 2\pi - \cos^{-1}) \\
 & & \cos \theta &= \vec{e} \cdot \vec{r}/er & \text{(if } v_r < 0 \text{ then } \theta = 2\pi - \cos^{-1})
 \end{aligned}$$

Oblateness of the Earth

$$\begin{aligned}
 a &= 6,378,137.0 \text{ m} & f^{-1} &= 298.257 & J_2 &= 1.0826 \times 10^{-3} \\
 \dot{\Omega} &= -\Upsilon \cos i & \dot{\omega} &= -\Upsilon \left(\frac{5}{2} \sin^2 i - 2 \right) & \Upsilon &= \frac{3}{2} \frac{\sqrt{\mu} J_2 R^2}{(1 - e^2)^2 a^{7/2}}
 \end{aligned}$$

Geocentric equatorial frame: $\vec{r} = r \cos \delta (\cos \alpha \hat{I} + \sin \alpha \hat{J}) + r \sin \delta \hat{K}$, where α and δ are the right ascension and declination.

Topocentric horizon frame: $\vec{\rho} = \rho \cos a (\sin A \hat{i} + \cos A \hat{j}) + \rho \sin a \hat{k}$, where a and A are the altitude and azimuth.

Topocentric equatorial frame: θ and ϕ are the sidereal time ($\theta = \theta_G + \Lambda$) and *geodetic* latitude, and α and δ are the *topocentric* right ascension and declination.

$$\begin{aligned}
 \hat{i} &= -\sin \theta \hat{I} + \cos \theta \hat{J} & \hat{j} &= -\sin \phi (\cos \theta \hat{I} + \sin \theta \hat{J}) + \cos \phi \hat{K} & \hat{k} &= \cos \phi (\cos \theta \hat{I} + \sin \theta \hat{J}) + \sin \phi \hat{K} \\
 \cos(\theta - \alpha) &= \left(\frac{\cos \phi \sin a - \sin \phi \cos A \cos a}{\cos \delta} \right) & \text{(if } \sin A > 0 \text{ then } (\theta - \alpha) = 2\pi - \cos^{-1}) \\
 \sin \delta &= \cos \phi \cos A \cos a + \sin \phi \sin a & \vec{\rho} &= \rho \cos \delta (\cos \alpha \hat{I} + \sin \alpha \hat{J}) + \rho \sin \delta \hat{K} \\
 \vec{r} &= \vec{R} + \vec{\rho} & \vec{R} &= R_C \cos \phi (\cos \theta \hat{I} + \sin \theta \hat{J}) + R_S \sin \phi \hat{K} \\
 R_S &= (1 - f)^2 R_\phi + H & R_C &= R_\phi + H & R_\phi &= \frac{R_e}{\sqrt{1 - (2f - f^2) \sin^2 \phi}}
 \end{aligned}$$

Orbital maneuvers: Hohmann transfer, bi-elliptic is more efficient if $r_C > 16r_A$.

Phasing maneuver (satellite paradox).

Plane change: $\Delta v = \sqrt{(v_{r1} - v_{r2})^2 + v_{\perp 1}^2 + v_{\perp 2}^2 - 2v_{\perp 1}v_{\perp 2} \cos \delta}$, most efficient if Δv_r is small and $v_{\perp 1}$ is small.

At periapse/apoapse: $\Delta v = \sqrt{(v_2 - v_1)^2 + 4v_1v_2 \sin^2(\delta/2)} \rightarrow 2v \sin(\delta/2)$ for a pure rotation).

Launch azimuth: $\cos i = \cos \phi \sin A$.

Orbit determination from the topocentric horizon frame

$$\vec{\Omega} = \omega_E \hat{K} \quad \dot{\vec{R}} = \vec{\Omega} \times \vec{R}$$

$$\dot{\delta} = \frac{-\dot{A} \cos \phi \sin A \cos a + \dot{a}(\sin \phi \cos a - \cos \phi \cos A \sin a)}{\cos \delta}$$

$$\dot{\alpha} = \omega_E + \frac{\dot{A} \cos A \cos a - \dot{a} \sin A \sin a + \dot{\delta} \sin A \cos a \tan \delta}{\cos \phi \sin a - \sin \phi \cos A \cos a}$$

$$\dot{\hat{\rho}} = (-\dot{\alpha} \sin \alpha \cos \delta - \dot{\delta} \cos \alpha \sin \delta) \hat{I} + (\dot{\alpha} \cos \alpha \cos \delta - \dot{\delta} \sin \alpha \sin \delta) \hat{J} + \dot{\delta} \cos \delta \hat{K} \quad \vec{v} = \dot{\vec{R}} + \dot{\rho} \hat{\rho} + \rho \dot{\hat{\rho}}$$

Plane change maneuvers

$$\cos i_2 = \cos \phi \sin A \quad i_2 \geq \phi \quad \Delta v = \sqrt{(v_{r2} - v_{r1})^2 + v_{\perp 1}^2 + v_{\perp 2}^2 - 2v_{\perp 1}v_{\perp 2} \cos \delta}$$

At apoapse or periapse, when $v_{r1} = v_{r2} = 0$, then $\Delta v = \sqrt{v_1^2 + v_2^2 - 2v_1v_2 \cos \delta}$

For a pure rotation, when $v_2 = v_1 = v$, then $\Delta v = 2v \sin\left(\frac{\delta}{2}\right)$

Relative motion/Tidal acceleration

$$\delta \ddot{\vec{r}} = -\frac{\mu}{R^3} \left[\delta \vec{r} - \frac{3\vec{R}}{R^2} \vec{R} \cdot \delta \vec{r} \right] \quad \ddot{\vec{R}} = -\mu \frac{\vec{R}}{R^3}$$

Clohessy-Wiltshire

$$\delta \ddot{x} - 3n^2 \delta x - 2n \delta \dot{y} = 0$$

$$\delta \ddot{y} + 2n \delta \dot{x} = 0$$

$$\delta \ddot{z} + n^2 \delta z = 0$$